
(To appear in **Optical Engineering**)

An Automated Stokesmetric Imaging Laser Radar System

Xue Liu¹, Shih Tseng¹, Renu Tripathi^{1,3} and M.S. Shahriar^{1,2,*}

¹*Department of Electrical Engineering and Computer Science*

²*Department of Physics and Astronomy*

Northwestern University, 2145 Sheridan Road, Evanston, IL 60208

³*Department of Physics & Pre-Engineering, Delaware State University, Dover, DE 19901*

**Corresponding author: shahriar@northwestern.edu*

Abstract: We report the design and implementation of a high-speed, automated laser radar (LADAR) system with sensitivity enhanced by polarimetric imaging technique. This LADAR is able to analyze the Stokes vector of the reflected light from a target at video-rate. With a polarization state generator and a polarization state analyzer, the system is capable of performing a complete Mueller matrix imaging of the scene under observation. This polarization-sensitive LADAR (PLADAR) is applied to various scenes and found to yield the ability to detect information which is indiscernible to a conventional, intensity-based LADAR.

Subject terms: Laser radar; Polarization imaging; Stokesmetry.

1. Introduction

Laser radar, also referred to as LADAR or LIDAR, uses eye-safe lasers to measure the range, and other properties of the targets. The LADAR shares many of the basic features of common microwave radars. But due to its extremely short operating wavelength and small beam size, it has been utilized extensively for defensive purposes such as in-port threat defense and shipboard protection systems since its invention [1-3].

The capabilities of a LADAR imaging system can be further enhanced if the light reflected from objects can be analyzed polarimetrically [4-7]. The term polarization, used to describe the complex direction of the electric field vector, plays an essential role in the interaction of light and matter [8]. Polarimetric imaging technique takes advantage of the fact that a given object emits and scatters light in a unique way depending on its properties, which allows us to distinguish objects with similar reflectivity but different polarimetric features. By breaking down the light into independent polarization components, one can often reveal occluded surface information such as the composition, texture and roughness. As such, polarimetric imaging is utilized widely in many situations. For example, astrophysicists use polarimeters to measure the spatial distribution of magnetic fields on the surface of the sun [9]. In the field of medical imaging, researchers analyze the polarization of light through coherence interferometry to yield the details about the optical properties of biological tissues. [10-12]. For defense applications, polarization imaging is often used to detect objects through smoke and obscurants such as vegetation or camouflage [13].

An arbitrary polarization state of light can be represented by a column vector called the Stokes vector, denoted as $\bar{S} \equiv (I \ Q \ U \ V)^T$. Each element in this vector is a linear combination of two orthogonal polarization components, which is illustrated explicitly in section 2. Over the years, many designs have been used to measure the Stokes vector of light such as a three cameras

architecture designed by Prosch [14], a liquid crystal polarization camera by Andreou and Wolff [15] and a dual piezoelastic modulator system by Stenflo and Povel [16].

In this paper, we report a fully automated polarization-sensitive LADAR (PLADAR) system that operates at a speed much higher than the classical Stokesmeter. This work builds on previous polarimetric studies carried out by our group [12,17,18]. Our system is capable of detecting all four components of the Stokes vector of light reflected from the object at video-rate. In addition, by varying the input polarization state, full Mueller matrix imaging is also realizable with such a system. An example of such full Muller-metric imaging employing variable input polarization was recently demonstrated by us in reference [12] in a somewhat different setting.

2. Stokes–Mueller Formalism

Polarization dependent reflectivity is quite common in nature. In fact, every surface alters the polarization of the incident radiation, if very slightly. Any arbitrary polarization state of the light can be represented by the well-known Stokes Vector, which is defined in terms of two orthogonal components of the oscillating electric field (denoted as E_x and E_y in this paper) perpendicular to the direction of propagation.

$$\bar{S}(\tau) \equiv \begin{pmatrix} I \\ Q \\ U \\ V \end{pmatrix} \equiv \begin{pmatrix} \langle E_x E_x^* \rangle + \langle E_y E_y^* \rangle \\ \langle E_x E_x^* \rangle - \langle E_y E_y^* \rangle \\ \langle E_x E_y^* \rangle + \langle E_y E_x^* \rangle \\ i \langle E_x E_y^* \rangle - i \langle E_y E_x^* \rangle \end{pmatrix} \quad (1)$$

where the brackets imply averaging over the observation time τ . I is the overall intensity, Q denotes the intensity difference between vertical and horizontal linear polarizations, U stands for the intensity difference between linear polarizations at $+45^\circ$ and -45° , and V is the intensity

difference between left and right circular polarizations. $I^2 = Q^2 + U^2 + V^2$ for completely polarized light and $I^2 > Q^2 + U^2 + V^2$ for partially polarized/unpolarized light,

In the Mueller–Stokes formalism, a 4×4 matrix, called the Mueller matrix (denoted as M), is often used to describe mathematically the transformation of the polarization of light by a medium [19]. The Mueller matrix relates the Stokes vectors of incoming and outgoing light (denoted as \vec{S}_{in} and \vec{S}_{out} , respectively) via the relation $\vec{S}_{out} = M\vec{S}_{in}$. The net Mueller matrix of any optical system is given by the ordered product of the Mueller matrices of the individual elements. From the elements of the Mueller matrix, we can infer the properties of the medium such as the reflectance, absorption and birefringence.

The Stokes–Mueller formalism is more comprehensive than the Jones vector approach, even though they both rely on linear algebra and matrix formalisms. Specifically, the Stokes vector encompasses any polarization state of light, whether it is completely or partially polarized, while the Jones vector deals only with fully polarized light.

3. Automated Inline Stokesmeter

A typical Stokesmetric imaging system consists of a Polarization State Generator (PSG) and a Polarization State Analyzer (PSA). The PSG controls the input polarization state of the light. The PSA analyzes the polarization of the light reflected by the object. The conventional PSA consists of a quarter wave plate (QWP) and a linear polarizer (LP), and requires the QWP to be inserted and removed between various readings [20]. Consequently, the speed of the system is greatly limited, which hinders it from being integrated into a real-time LADAR system. We proposed and demonstrated a PSG consisting of a pair of Liquid Crystal Retarders (LCR's) and a polarizer

in series to circumvent this problem in our previous study [17]. Due to the anisotropic optical property, an LCR can produce arbitrary retardation based on the voltage applied on it. This flexibility allows a configuration of the PSA free of insertion/removal, as shown in Figure 1. This architecture, which has been investigated theoretically as well as experimentally [21-24], forms the foundation of our PLADAR system. Like the classical Stokesmeter, four measurements are required to determine the Stokes vector of the unknown light. Table 1 lists the required phase delays of the LCR's for each measurement for such an architecture.

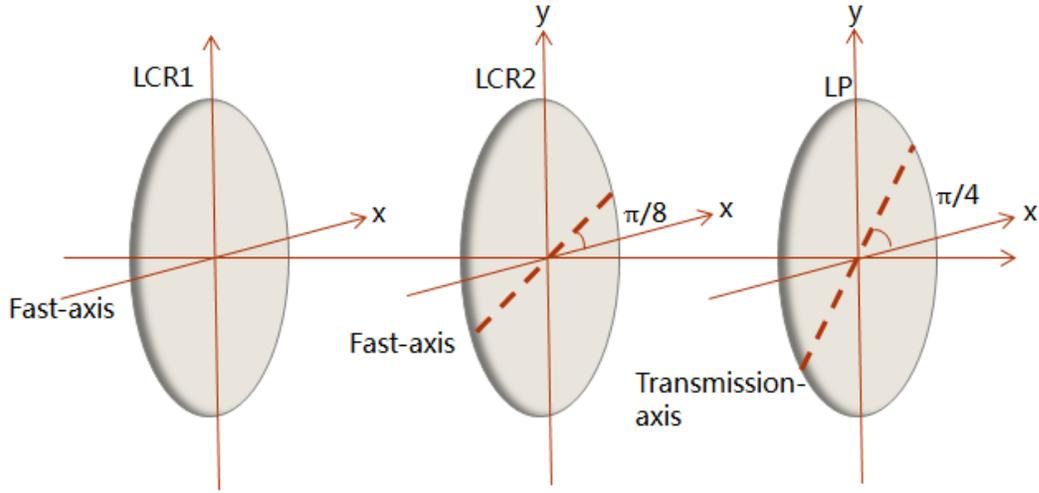


Figure 1: Automated inline Stokesmeter based on the Liquid Crystal Retarders.

Table 1: Phase delay of the LCR's for each measurement of the inline Stokesmeter

Detected Intensity	LCR1	LCR2
I_1	$\pi/2$	0
I_2	π	0

I_3	0	π
I_4	0	0

The Stokes vector can be calculated based on the four distinctly measured intensities with the following equation:

$$\begin{bmatrix} I_1 \\ I_2 \\ I_3 \\ I_4 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 1 \\ 1 & 0 & -1 & 0 \\ 1 & 1 & 0 & 0 \\ 1 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} I \\ Q \\ U \\ V \end{bmatrix} \quad (2)$$

4. Fully Automated Polarimetric Ladar System

The basic configuration for the PLADAR is illustrated schematically in Figure 2. A pulsed laser at 1571 nm (pulse width = 6.5 ns) is used as the light source. The Stokes imaging occurs in the following order: At the pre-detection stage, the computer sends signals to the LCR controller as

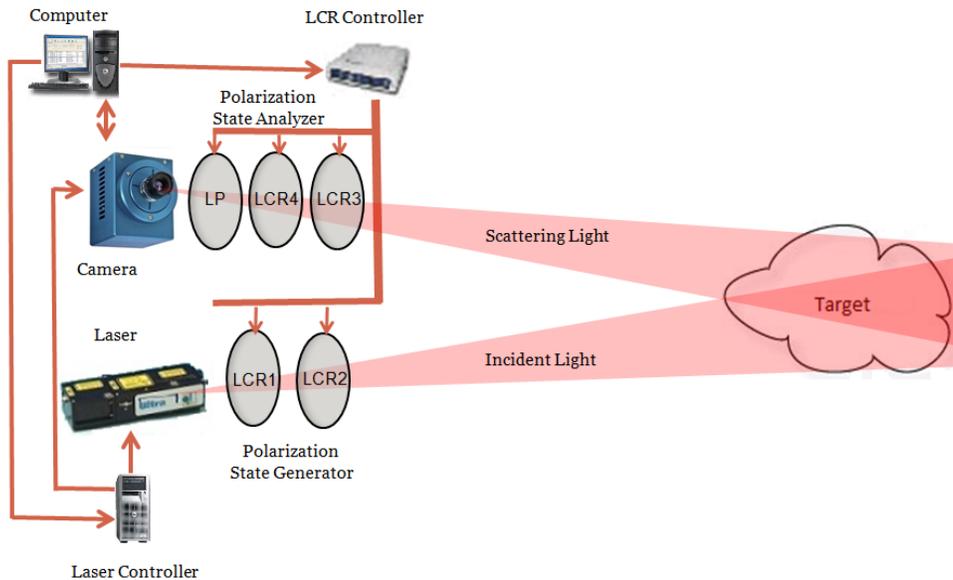


Figure 2: Schematic illustration of the fully-automated imaging Polarimetric LADAR.

well as the laser controller. The LCR controller sets the LCR's in proper states. Once the LCR's are ready, the laser controller triggers the pulse and the camera simultaneously. The source laser is first sent through the PSG to produce the desired polarization before hitting the target. The back-scattered light transmits the PSA and is then detected by the infrared camera, which has a CCD array of 300×256 pixels with $30 \times 30 \mu m^2$ pixel size. An exposure time window is applied to the camera to eliminate the unwanted background noise. The recorded intensity data array is transmitted back to the computer where the Stokes vector of each pixel is calculated simultaneously by a computer program, thus producing four distinct images, corresponding to the four Stokes parameters, for the input scene.

The operation speed of this PLADAR system is determined primarily by the response time of the LCR's. As mentioned in section 3, to analyze the Stokes vector, four independent measurements with different LCR phase delays are required. For the LCR's used in our experimental setup, it takes 15 ms for the retardation to increase from 0 to π , and 5 ms to relax from π to 0. The activation and relaxation of phase delay of each LCR in a measurement cycle is shown in Figure 3. Compared with the response time of the LCR's, the exposure window for light detection (1 μs for the camera in our setup) is negligible. The intensities I_1 , I_2 , I_3 and I_4 listed in Table-1 are recorded at times T_1 , T_2 , T_3 and T_4 , respectively. For our current setup, completing one cycle of 4 measurements for any input scene takes about 35ms, enabling the LADAR system to operate at video-rate.

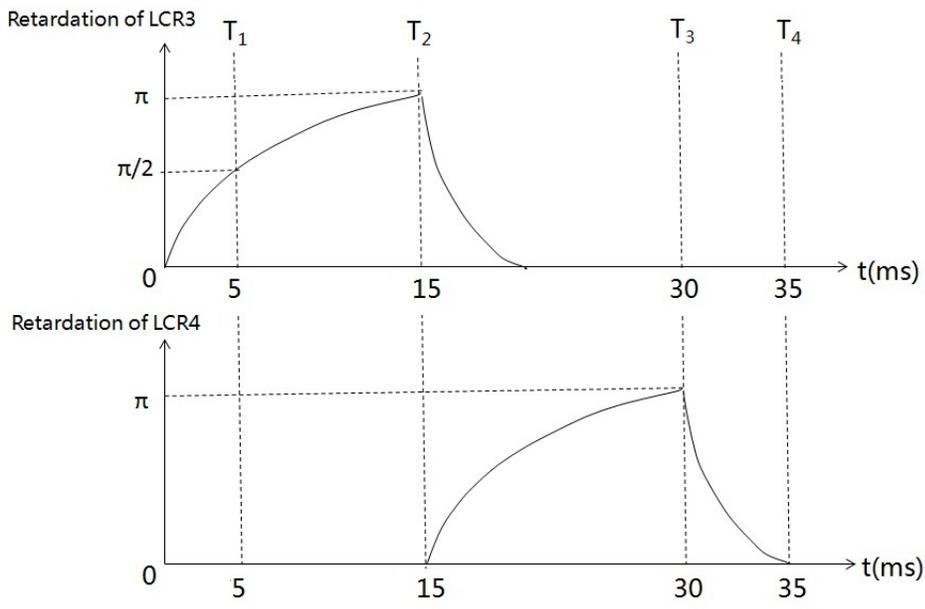


Figure 3: The activation and relaxation of phase delay of each LCR in a measurement cycle.

5. Stokesmetric Imaging of Various Scenes

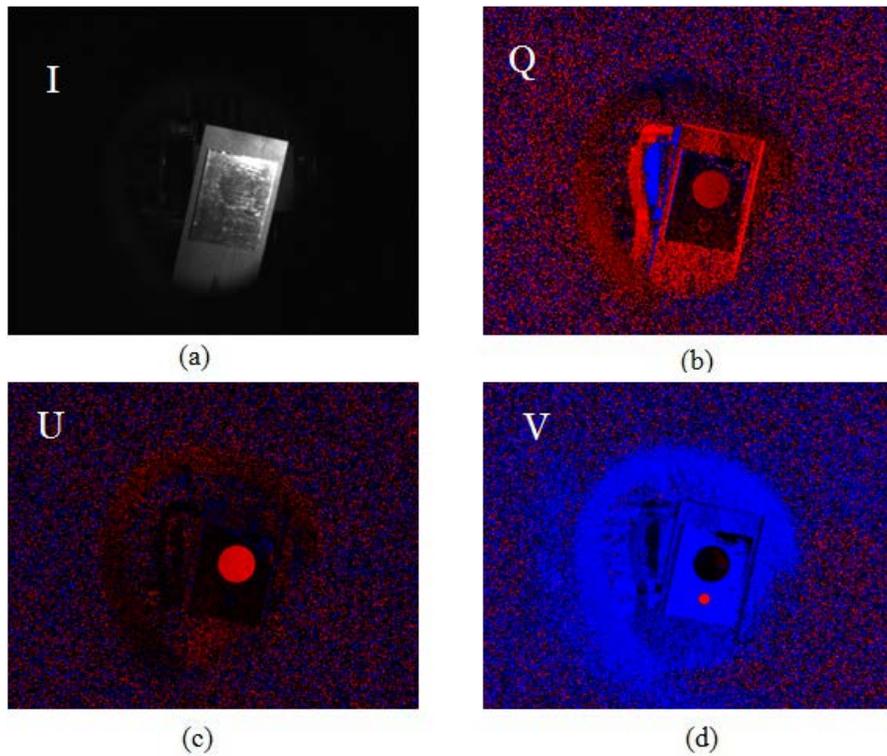


Figure 4: Stokesmetric images of objects illuminated by right circularly polarized light.

We first tested the PLADAR imaging system with an artificial scene comprising of optical elements with well-defined polarimetric properties. Specifically, we mounted a set of polarizer (at 45°) and a QWP on a piece of metal. Right circularly polarized light is generated using the PSG to illuminate this artificial scene. The resultant four Stokes vector images of the reflected light are shown in figure-4. The false red/blue color is imposed by the software to represent positive/negative value, respectively. Image (a), corresponding to the I element in the Stokes vector, contains only intensity-based information. Figures 'b' through 'd' are the Q, U and V images of the Stokes vector. Since our target objects are almost transparent, they can barely be distinguished from the background. However, because the tilted polarizer only selects 45° polarization to pass while the rest of the scene does not, it is easily distinct in the picture (c) containing only the U element. Similarly, because the QWP produces an extra π phase shift to the incident light in a round trip, the V image of it appears red, while the surroundings appear blue. Such an obvious contrast makes it possible to detect easily the almost transparent QWP in (a).

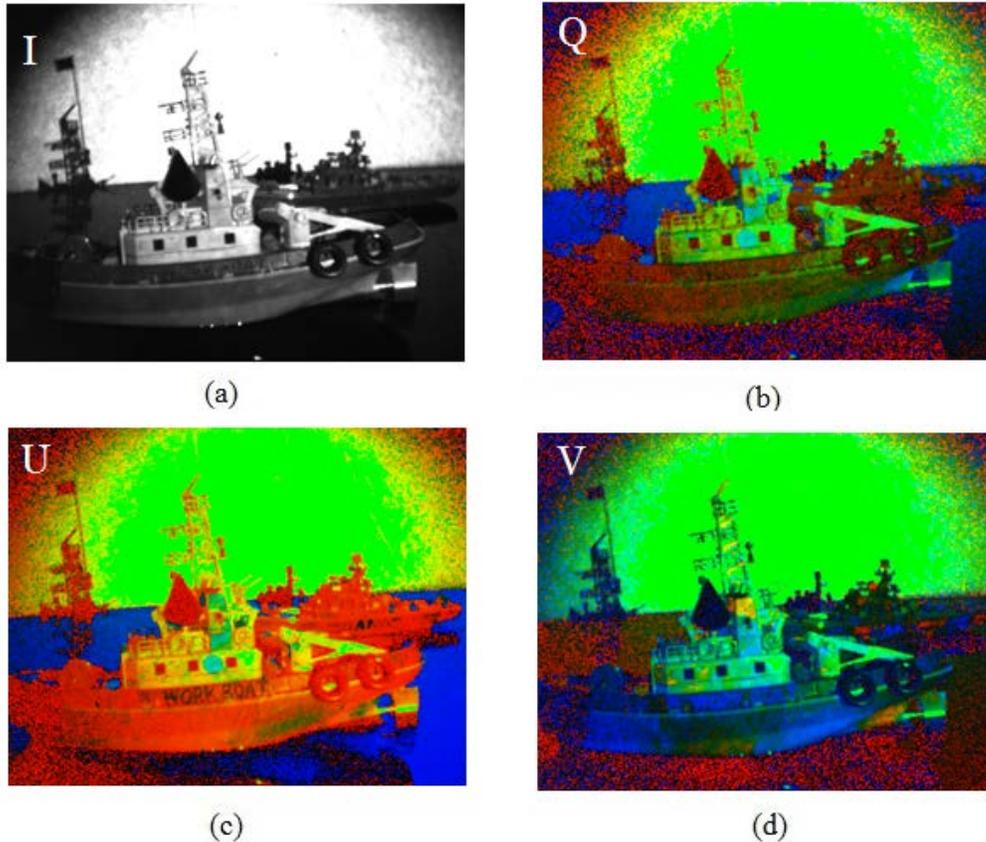


Figure 5: Stokesmetric images produced under right circularly polarized light, for an artificial harbor.

After validating the capabilities of the automated PLADAR system, we used it to capture the Stokesmetric images of targets merged in more complicated scenes. Figure 5 shows a sample image set obtained for one such scene. We imitated the scene of a harbor by placing model boats on still water. A crystal is placed between two windows of the boat and a transparent tape is attached to the body of the boat to form the phrase “work boat”. As can be seen, both the crystal and the letters can hardly be recognized from Figure 5 (a). However, image (c), corresponding to the U vector (containing purely 45° polarized light) clearly depicts the crystal due to its birefringence. The word “work boat” is also readable from image (c) from the body of the boat because of its different polarimetric signature from the surroundings. This ability to identify and

distinguish various surfaces with a very high contrast makes the PLADAR a promising tool for surveillance and tracking applications. Note that for this imaging system, the angle of view (AOV) is primarily determined by the dimension of the CCD sensor and the effective focal length (EFL) of the camera. In our current setup, the $9 \times 7.7 \text{mm}^2$ CCD sensor area and 16mm EFL of the relay lenses yield an AOV of 31.5° and 27.0° in the horizontal and vertical directions, respectively. The AOV can be further enhanced by increasing the size of the CCD for applications such as the wide-area surveillance.

By measuring the Stokes parameter images for a set of different input polarizations states, it would be possible to produce the full set of Muellermetric images, in a manner similar to what we have demonstrated recently for a polarimetric optical coherence tomography system [12]. This work is in progress, and will be reported in a future publication.

Summary

The polarimetric information embedded in the images of a scene is not detected by an intensity-image based LADAR system. In this paper, we report the implementation of a fully automated, video-rate polarimetric LADAR system that is capable of detecting the Stokes vector of the reflected light. The polarimetric images of the artificial scenes reveal information not detectable by the intensity-based LADAR. It is also possible to perform full Mueller-metric imaging with such a system by controlling the input polarization state of light.

Acknowledgements

This work is supported in part by AFOSR grant #FA9550-06-1-0466 and #FA9550-10-1-0228, NASA Grant #NNX09AU90A, and NSF Crest grant #0630388. One of the authors, Renu

Tripathi, acknowledges support from National Science Foundation NSF-CREST (grant # 0630388), and NASA URC-5 (grant # NNX09AU90A).

References

- [1]. L.J. Sullivan, "Infrared Coherent Radar," SPIE **227**, Development of Coherent Laser Radar at Lincoln Laboratory, 148–161. (1980).
- [2]. Walter G. Egan, *Optical Remote Sensing Science and Technology*, (Marcel Dekker, 2004).
- [3]. Henderson, B.W. " 'Firefly' Laser Experiment Successful in Measuring Inflatable Decoy Motion," Aviat. Week Space Technol. **132**, 75 (1990).
- [4]. T. Nee and S. F. Nee, "Infrared polarization signatures for targets," Proc. SPIE **2469**, 231–241 (1995).
- [5]. P.J. Curran, "Polarized visible light as an aid to vegetation classification," Remote Sens. Environ. **12**, 491-499, (1982).
- [6]. M. J. Duggin, "Imaging polarimetry in scene element discrimination," Proc. SPIE **3754**, 108-117 (1999).
- [7]. Shih-Yau Lu and Russell A. Chipman, "Interpretation of Mueller matrices based on polar decomposition," J. Opt. Soc. Am. A **13**, 1106-1113 (1996)
- [8]. Brian J. DeBoo, Jose M. Sasian, and Russell A. Chipman, "Depolarization of diffusely reflecting man-made objects," Appl. Opt. **44**, 5434-5445 (2005)
- [9]. Haosheng Lin and Thomas Rimmele, "The Granular Magnetic Fields of the Quiet Sun", Astr. Phys. J. **514**, 448-455 (1999)

-
- [10]. Johannes F. de Boer, Thomas E. Milner, and J. Stuart Nelson, "Determination of the depth-resolved Stokes parameters of light backscattered from turbid media by use of polarization-sensitive optical coherence tomography," *Opt. Lett.* **24**, 300-302 (1999)
- [11]. Erich Götzinger, Michael Pircher, Markus Sticker, Adolf F. Fercher and Christoph K. Hitzenberger, "Measurement and imaging of birefringent properties of the human cornea with phase-resolved, polarization-sensitive optical coherence tomography", *J. Biomed. Opt.* **9**, 94-102 (2004)
- [12]. Xue Liu, Shih C. Tseng, Renu Tripathi, Alexander Heifetz, Subramanian Krishnamurthy, M.S. Shahriar, "White light interferometric detection of unpolarized light for complete Stokesmetric optical coherence tomography," *Opt. Comm.* **284**, 3497–3503 (2011)
- [13]. Duggin, M.J.; Kinn, G.J.; Schrader, M. "Enhancement of vegetation mapping using Stokes parameter images," *Proc. SPIE* **3121**, 307–313 (1997).
- [14]. T. Prosch, D. Hennings, and E. Raschke, "Video polarimetry: a new imaging technique in atmospheric science," *Appl. Opt.* **22**, 1360-1363 (1983)
- [15]. L. Wolff and A. Andreou, "Polarization camera sensors," *Image Vision Comput.* **13**, 497–510 (1995).
- [16]. Hanspeter Povel, Hans Aebersold, and Jan O. Stenflo, "Charge-coupled device image sensor as a demodulator in a 2-D polarimeter with a piezoelastic modulator," *Appl. Opt.* **29**, 1186-1190 (1990).
- [17]. X. Liu, A. Heifetz, S.C. Tseng, and M.S. Shahriar, "High-speed inline holographic Stokesmeter imaging," *Appl. Opt.* **48**, 3803-3808 (2009).
- [18]. Jong-Kwon Lee, Xue Liu, and M. S. Shahriar, "Volume-grating Stokesmeter based on photonic bandgap structures," *Appl. Opt.* **48**, 3212-3215 (2009).

-
- [19]. D. Goldstein, *Polarized Light*, 3rd ed. CRC Press (2010).
- [20]. C. Brosseau, *Fundamentals of Polarized Light: A Statistical Optics Approach*, Wiley, New York (1988).
- [21]. J. M. Bueno, “Polarimetry using liquid-crystal variable retarders: theory and calibration”, *J. Opt. A.* **2**, 216–222 (2000).
- [22]. De Martino A, Garcia-Caurel E, Laude B and Drévillon B, “Optimized Mueller polarimeter with liquid crystals”, *Opt. Lett.* **28**, 616–618 (2003).
- [23]. B. Laude-Boulesteix, A. De Martino, B. Drévillon, and L. Schwartz, “Mueller polarimetric imaging system with liquid crystals”, *Appl. Opt.* **43**, 2824-2832 (2004).
- [24]. Goudail F, Terrier P, Takakura Y, Bigu´e L, Galland F and Devlaminck V, “Target detection with a liquidcrystal-based passive Stokes polarimeter,” *Appl. Opt.* **43** 274–282 (2004).